

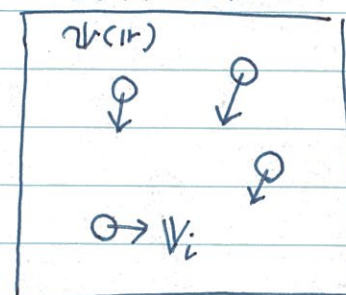
The effects of hydrodynamics^{No. 3}

流体力学相互作用の取り扱い.

1). Colloidal Dispersions

$$V_m = \sum_m H_{mm} \cdot F_m$$

溶媒の速度場.



i). Very Dilute Solution....

$$H_{mm} = \frac{\eta \delta_{mm}}{\zeta}$$

$$\rightarrow V_m^\alpha = F_m^\alpha / \zeta \quad \text{互に独立}$$

ii) Oseen tensor

$$V_m = v(R_m), \quad g(r) = \sum_m F_m \delta(r - R_m)$$

$$\begin{cases} \textcircled{1} & \nabla \cdot v(r) = 0 \quad (\text{連続の式}) \\ \textcircled{2} & \eta \nabla^2 v(r) + \nabla P = -g(r) \quad (\text{Navier-Stokes 式}) \end{cases}$$

$$\text{フーリエ変換} \quad v_k \equiv \frac{1}{V} \int d^3r v(r) \exp(i k \cdot r)$$

$$\begin{cases} \textcircled{3} & k \cdot v_k = 0 \\ \textcircled{4} & -\eta k^2 v_k - i k P_k = -g_k \end{cases}$$

$$\mathbf{k} \cdot \textcircled{4} = \textcircled{3} \epsilon \mathbf{f}' \lambda$$

$$-\eta k^2 \mathbf{k} \cdot \psi_{\mathbf{k}} - i k^2 P_{\mathbf{k}} = -\mathbf{k} \cdot \mathcal{J}_{\mathbf{k}}$$

$$\textcircled{5} \quad P_{\mathbf{k}} = \frac{-i}{k^2} \mathbf{k} \cdot \mathcal{J}_{\mathbf{k}}$$

$$\textcircled{5} \epsilon \textcircled{4} = \mathbf{f}' \lambda$$

$$-\eta k^2 \psi_{\mathbf{k}} - i k \frac{-i k}{k^2} \cdot \mathcal{J}_{\mathbf{k}} = -\mathcal{J}_{\mathbf{k}}$$

$$\psi_{\mathbf{k}} = \frac{1}{\eta k^2} \left(\mathbb{I} - \frac{\mathbf{k} \mathbf{k}}{k^2} \right) \cdot \mathcal{J}_{\mathbf{k}}$$

$$= H_{\mathbf{k}} \cdot \mathcal{J}_{\mathbf{k}}$$

↓ 実空間に戻すと...

$$\psi(\mathbf{r}) = \int d\mathbf{r}' H(\mathbf{r} - \mathbf{r}') \cdot \mathcal{J}(\mathbf{r}')$$

$$H(\mathbf{r}) = \frac{1}{8\pi\eta r} \left(\mathbb{I} + \frac{\mathbf{r} \mathbf{r}}{r^2} \right)$$

Oseen tensor //

$$V_m = \psi(R_m)$$

$$= \int d\mathbf{r}' H(\mathbf{R}_m - \mathbf{r}') \cdot \mathcal{F}(\mathbf{r}')$$

$$= \int d\mathbf{r}' H(\mathbf{R}_m - \mathbf{r}') \cdot \sum_m \mathbf{F}_m \delta(\mathbf{r}' - \mathbf{R}_m)$$

$$= \sum_m H(\mathbf{R}_m - \mathbf{R}_m) \cdot \mathbf{F}_m$$

$$= \sum_m H_{mm} \cdot \mathbf{F}_m$$

Oseen tensor については、 H_{mm} が発散してしまう。
2つの

$$\therefore \begin{cases} H_{mm} = \frac{\mathbb{I}}{\zeta} \\ H_{mm} = \frac{1}{8\pi\eta r} \left(\mathbb{I} + \frac{\mathbf{r}\mathbf{r}}{r^2} \right) \quad m \neq m \end{cases}$$

とするとこれが正しい。

Oseen tensor を使う方法の問題点

i). 粒子の大きさ a を入れるのが困難。

ii). 粒子の距離が近しい場合は使えない。

↳ または Navier - Stokes を解く必要がある。

Unifying Concepts in
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$$V_m = \sum_m H_{mm} \cdot F_m$$

Oseen tensor

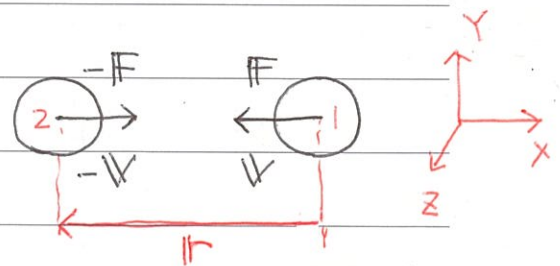
$$\begin{cases} H_{mm} = \frac{\text{I}}{6\pi\eta a} & r = |r_m - r_m| \\ H_{mm} = \frac{1}{8\pi\eta r} \left[\text{I} + \frac{r r}{r^2} \right] & m \neq m \end{cases}$$

RPY tensor

$$\begin{cases} H_{mm} = \frac{\text{I}}{6\pi\eta a} \\ H_{mm} = \frac{1}{8\pi\eta r} \left[\text{I} + \frac{r r}{r^2} + \frac{2}{3} \left(\frac{a}{r} \right)^2 \left(\text{I} - \frac{3r r}{r^2} \right) \right] \end{cases}$$

Quiz

i) squeeze



$$\begin{cases} V_1 = H_{11} \cdot F_1 + H_{12} \cdot F_2 \\ V_2 = H_{22} \cdot F_2 + H_{21} \cdot F_1 \end{cases}$$

$$\begin{aligned} F &= (F, 0, 0) \\ V &= (V, 0, 0) \\ r &= (r, 0, 0) \end{aligned}$$

Oseen

$$\begin{aligned} V &= \frac{1}{6\pi\eta a} F + \frac{1}{8\pi\eta r} [1 + 1] (-F) \\ &= \left[\frac{1}{6\pi\eta a} - \frac{1}{4\pi\eta r} \right] F \end{aligned}$$

$$\frac{V}{F} = \frac{1}{6\pi\eta a} - \frac{1}{4\pi\eta r}$$

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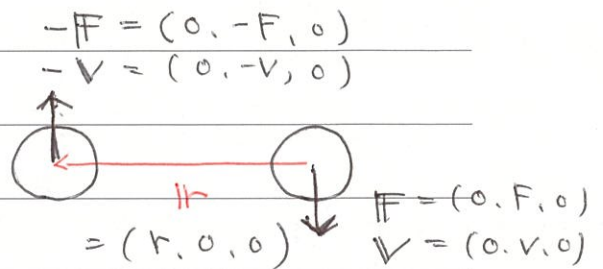
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RPY

$$\frac{V}{F} = \frac{1}{6\pi\eta a} - \frac{1}{4\pi\eta r} + \frac{1}{6\pi\eta r} \left(\frac{a}{r}\right)^2$$

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ii) tangential.



Oseen

$$V = \frac{1}{6\pi\eta a} F + \frac{1}{8\pi\eta r} [1 + 0] (-F)$$

$$= \left[\frac{1}{6\pi\eta a} - \frac{1}{8\pi\eta r} \right] F$$

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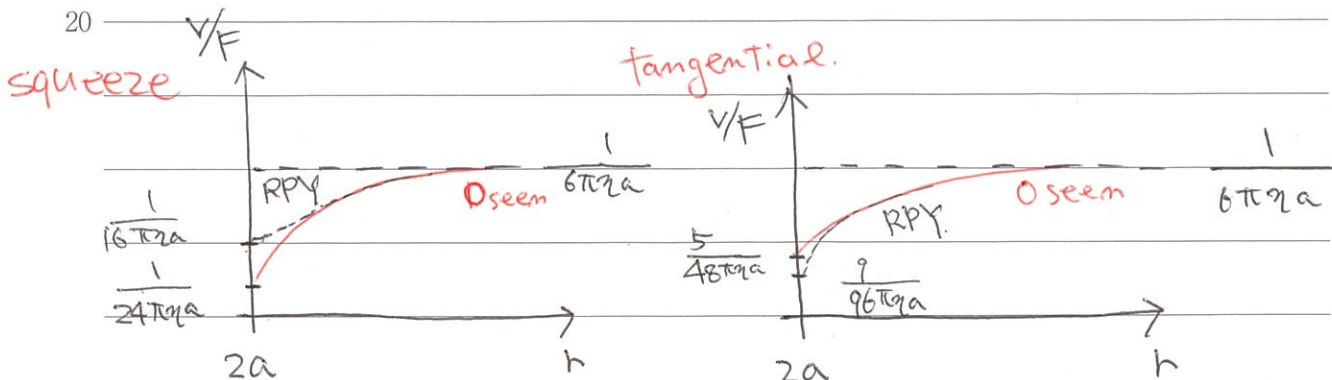
$$\frac{V}{F} = \frac{1}{6\pi\eta a} - \frac{1}{8\pi\eta r}$$

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RPY

$$\frac{V}{F} = \frac{1}{6\pi\eta a} - \frac{1}{8\pi\eta r} - \frac{1}{12\pi\eta r} \left(\frac{a}{r}\right)^2$$

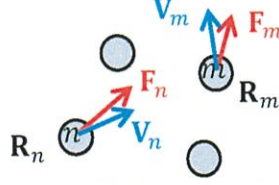
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3. Summary for hydrodynamic interactions (HI) between particles

Ryoichi Yamamoto

1. Mobility tensor



Suppose that a collection of N spherical particles, all having the same radius a , are suspended in an incompressible fluid with the viscosity η . Let $\mathbf{R}_1, \mathbf{R}_2, \dots, \mathbf{R}_N$ be the positions of the particles and $\mathbf{F}_1, \mathbf{F}_2, \dots, \mathbf{F}_N$ be the force acting on them. We assume that there are no external torques acting on the particles and inertia effects are all neglected due to very small Re . Then the velocities of the particles are written as

$$\mathbf{V}_n = \frac{d\mathbf{R}_n}{dt} = \sum_{m=1}^N \mathbf{H}_{nm} \cdot \mathbf{F}_m$$

by using the mobility tensor \mathbf{H}_{nm} . Three representations of \mathbf{H}_{nm} with different levels of approximations are summarized below.

I) No HI:

$$\begin{aligned} \mathbf{H}_{mm} &= \frac{1}{6\pi\eta a} \mathbf{I} \\ \mathbf{H}_{nm} &= 0 \quad (n \neq m) \end{aligned}$$

II) Oseen tensor:

$$\begin{aligned} \mathbf{H}_{mm} &= \frac{1}{6\pi\eta a} \mathbf{I} \\ \mathbf{H}_{nm} &= \frac{1}{8\pi\eta r} \left[\mathbf{I} + \frac{\mathbf{r}\mathbf{r}}{r^2} \right] \quad (n \neq m) \end{aligned}$$

III) Rotne-Prager-Yamakawa (RPY) tensor:

$$\begin{aligned} \mathbf{H}_{mm} &= \frac{1}{6\pi\eta a} \mathbf{I} \\ \mathbf{H}_{nm} &= \frac{1}{8\pi\eta r} \left[\mathbf{I} + \frac{\mathbf{r}\mathbf{r}}{r^2} + \frac{2}{3} \left(\frac{a}{r} \right)^2 \left(\mathbf{I} - \frac{3\mathbf{r}\mathbf{r}}{r^2} \right) \right] \quad (n \neq m) \end{aligned}$$

Here,

$$\mathbf{I} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \quad \mathbf{r} = \mathbf{R}_n - \mathbf{R}_m, \quad r = |\mathbf{r}|$$